

The total strain energy within any annular area of inner and outer radii  $r_0$  and  $R$  is

$$\Phi = \int_0^{2\pi} \int_{r_0}^R \frac{1}{2} \sigma_{ij} \epsilon_{ij} r dr d\theta \sim \int_0^{2\pi} \int_{r_0}^R r^{(2\lambda+1)} dr d\theta.$$

Because  $\Phi < \infty$  and  $\lambda > -1$ , the physically admissible values of  $\lambda$  are

$$\lambda = -\frac{1}{2}, 0, \frac{1}{2}, 1, \frac{3}{2}, 2, \dots, \frac{Z}{2},$$

where  $Z$  is an integer.

Taking  $\lambda = -1/2$

$$\chi = r^{3/2} A_1 \left[ \cos \frac{\theta}{2} + \frac{1}{3} \cos \frac{3\theta}{2} \right] + O(r^2) + O(r^{5/2}) + \dots,$$

$$\sigma_{ij} = A_1 r^{-1/2} \tilde{\sigma}_{ij}^I(\theta) + O_{ij}(r^0) + O_{ij}(r^{1/2}) + \dots$$

Rewriting

$$A_1 = K_I / \sqrt{2\pi r},$$

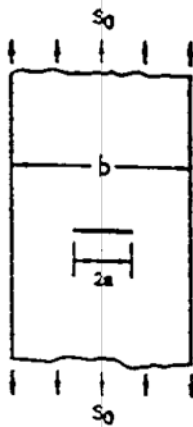
$$\sigma_{ij} = \frac{K_I}{\sqrt{2\pi r}} \tilde{\sigma}_{ij}^I(\theta) + T \delta_{ix} \delta_{jx}$$

+ (terms which vanish at crack tip),

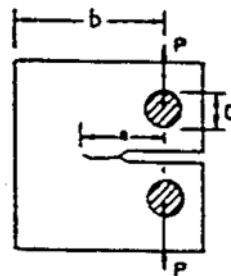
$K_I$  is the stress intensity factor for mode I.  $\delta_{ij}$  is the Kronecker delta.

# Possible Crack Configurations (I)

SS01

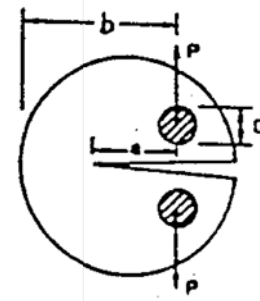


SS02



$B = THK$   
 $S_0 = P/BD$

SS03



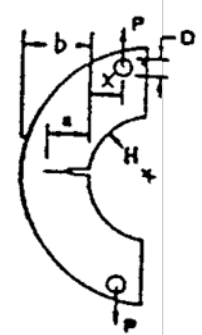
$B = THK$   
 $S_0 = P/BD$

## **Real Cracks in Real Structures**

- **Most structures and components contain surface cracks**
- **Surface cracks may be idealized as**
  - **quarter elliptic cracks/corner cracks**
  - **semi-elliptical/thumbnail cracks**
  - **elliptical/embedded cracks**
- **Cracks may be present at holes, pins or other stress concentrators**
- **Cracks may be subjected to combinations of bending, tension, torsion**
- **Cracks may also have multiple crack combinations**
- **Simple fracture mechanics approach applicable in many cases**

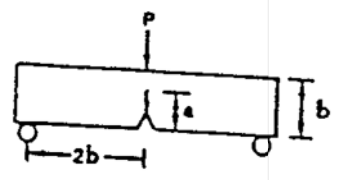
# Possible Crack Configurations (II)

3304



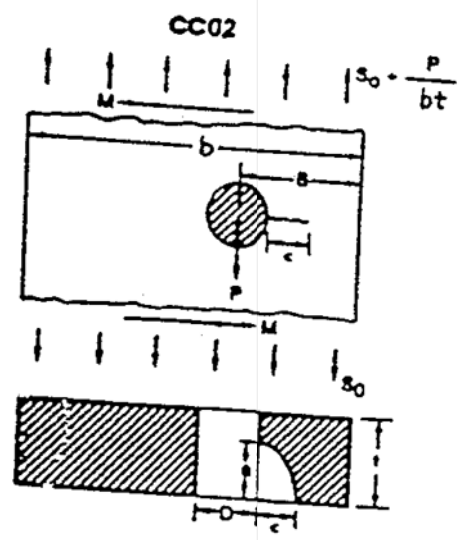
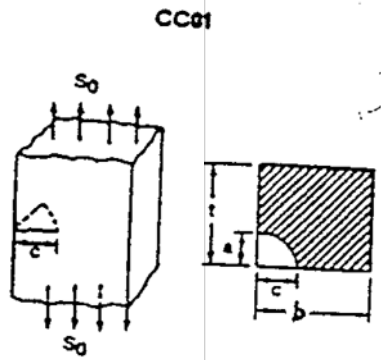
B = THK  
 $S_0 = P/BD$

3305

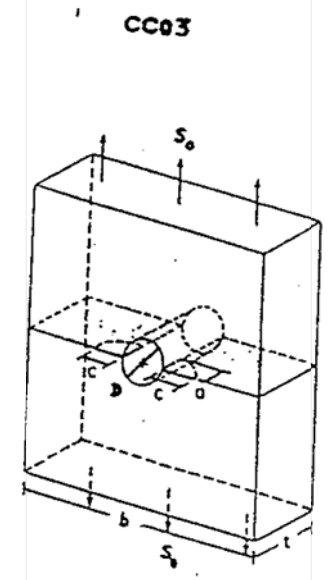


B = THK  
 $S_0 = \frac{6P}{Bb}$

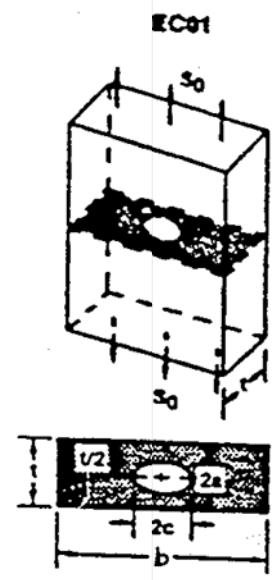
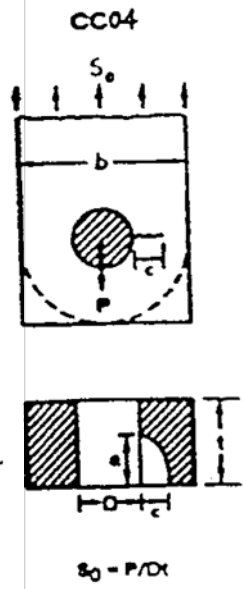
# Possible Crack Configurations (III)



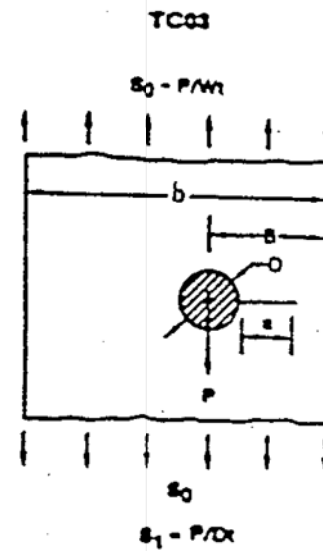
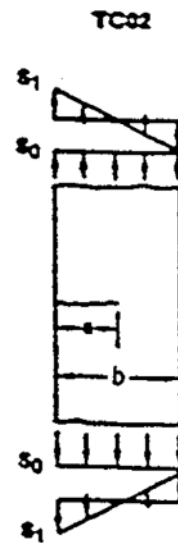
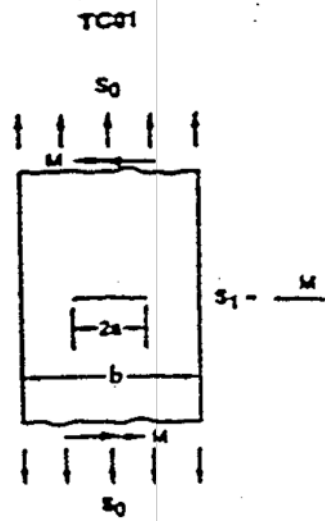
$$\sigma_1 = P/dt \quad \sigma_2 = \frac{6M}{bt^2}$$



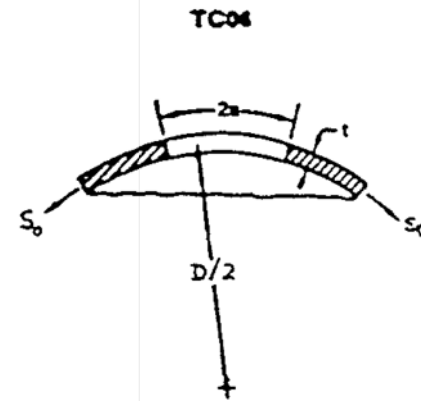
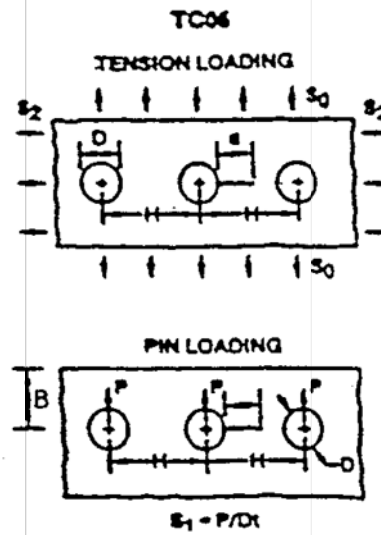
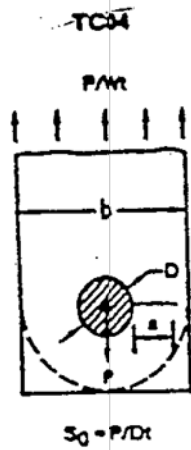
# Possible Crack Configurations (IV)



# Possible Crack Configurations (V)

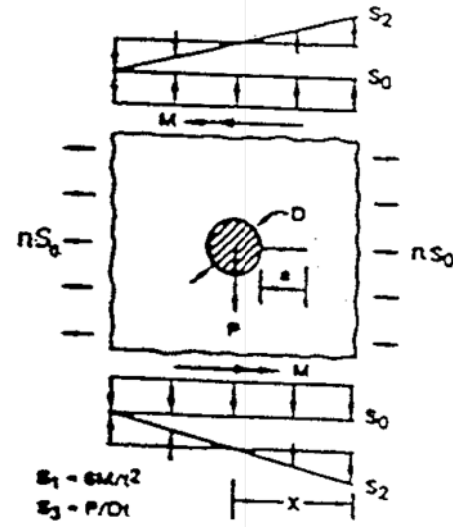
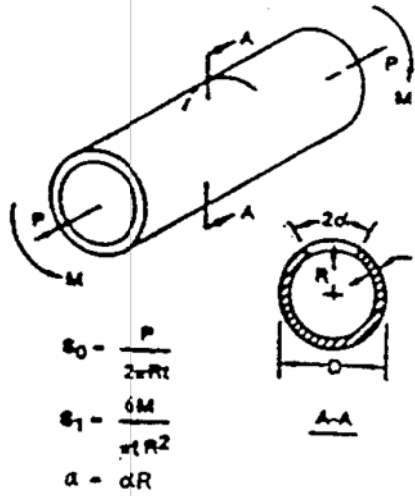
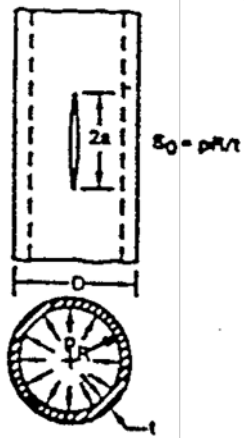


# Possible Crack Configurations (VI)

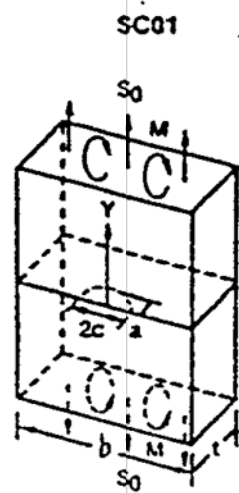




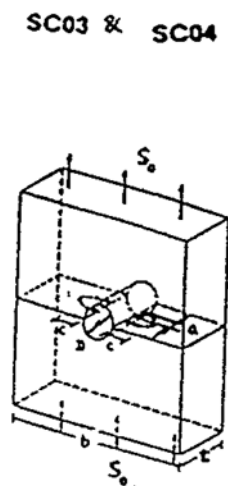
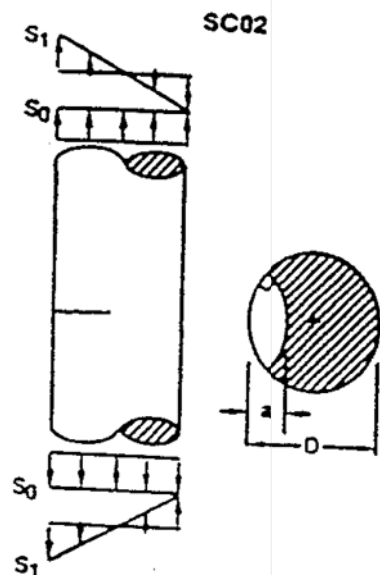
# Possible Crack Configurations (VII)



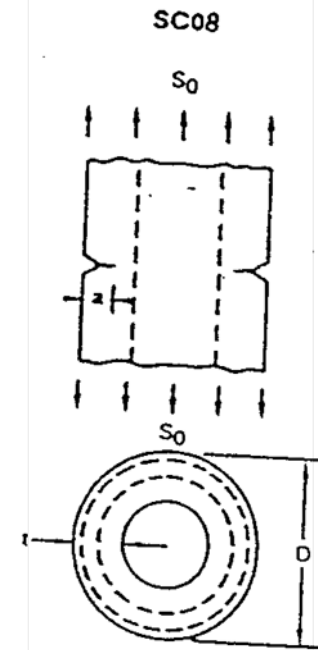
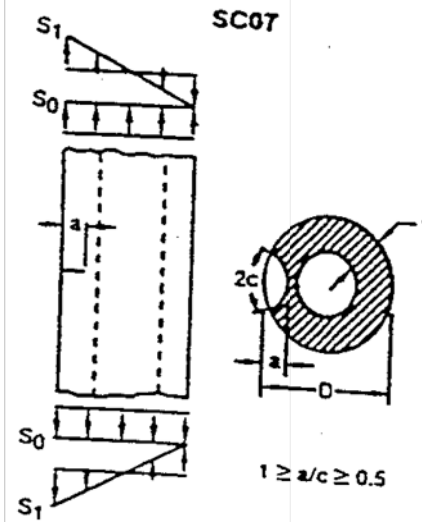
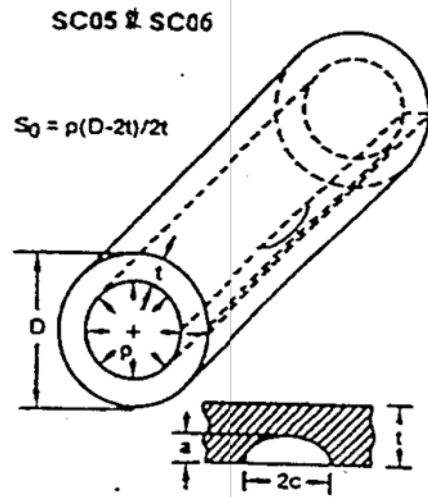
# Possible Crack Configurations (VIII)



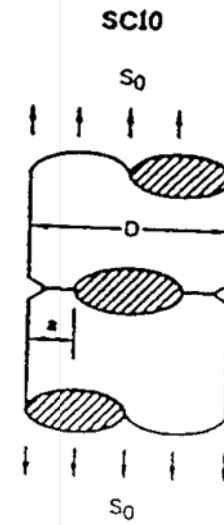
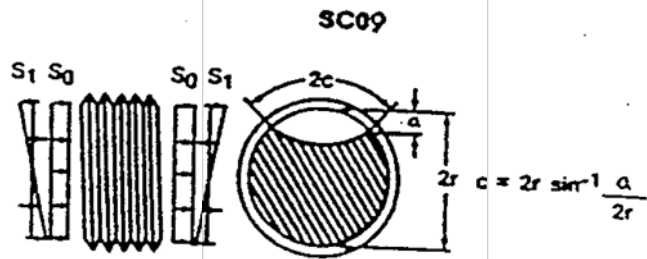
$$S_1 = \frac{6M}{wt^2}$$



# Possible Crack Configurations (IX)



# Possible Crack Configurations (X)



# Monotonic Plastic Zone Sizes

- Size depends on stress state
- Irwin (1966) showed that:
  - $PZS = (A/3 \cdot \pi) \cdot [DK/s_y]$
  - $A = 3$  (plane strain)
  - $A = 1$  (plane stress)
- Other expressions by Dugdale and Barrenblatt

# Cyclic Plastic Zone

- Smaller than monotonic plastic zone
- Affected by reversed plasticity
- $PZS = (1/\pi) * \{DK/2s_y\}$  (Rice, 1967)
- Residual wake left behind crack-tip
- Residual stresses induced as a result
- May promote crack closure

# Crack Opening Displacement

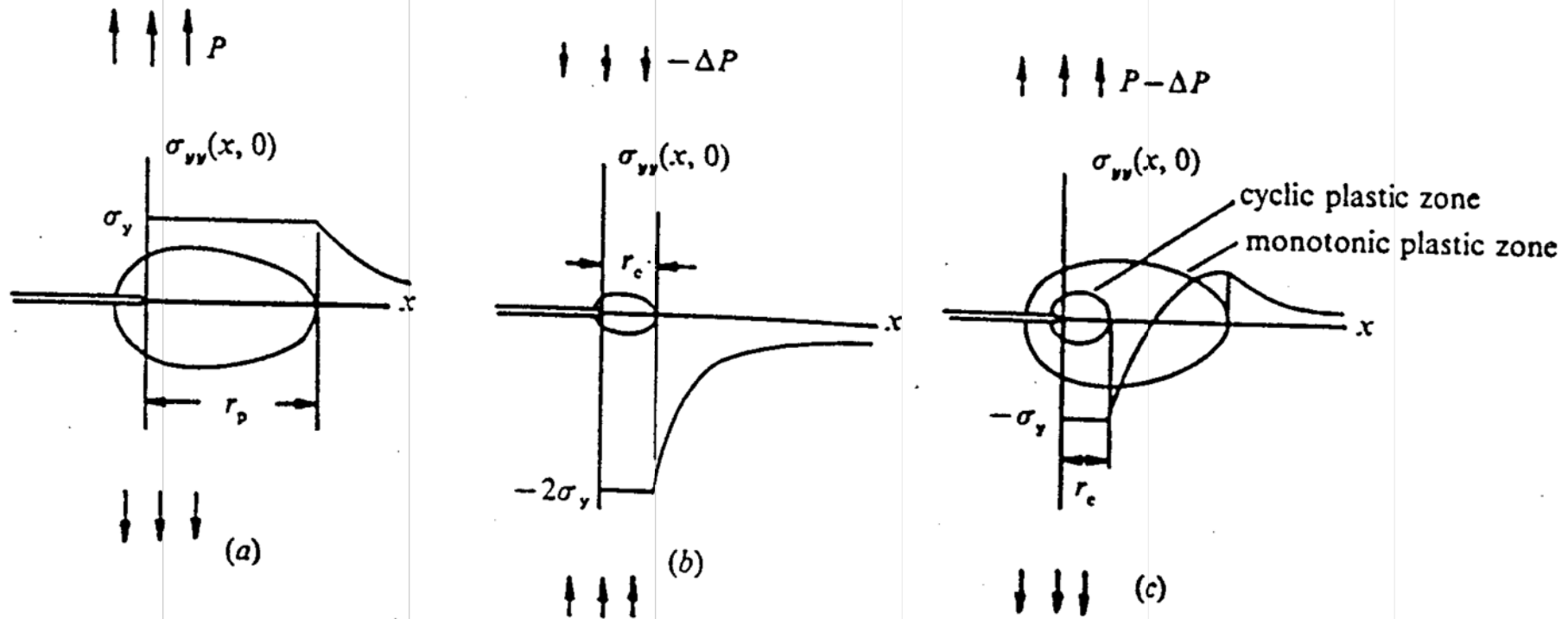
- Change in CTOD between max and min load
- Function of  $\Delta K$ ,  $E$  and yield stress
- Crack extension related to this change
- Crack opens to maximum load
- Extension due to plastic processes

# Plastic Zone Size Under Cyclic Loading

- PZ existence under cyclic loading long recognized
- Paris (1960), McClintock (1963) & Rice (1967)
- Experimental evidence by Hahn, Hoagland & Rosenfield (1972)
- PZ controlled by reversed plasticity
- PZ much smaller under cyclic loading



# Schematic representation of the development of cyclic plastic zone upon unloading (After Rice, 1967)



$$r_c \approx \frac{1}{\pi} \left( \frac{\Delta K_I}{2\sigma_y} \right)^2$$

- (a) Monotonic plastic zone created by a far-field load  $P$ .  
 (b) Stress distribution due to the reduction of the load by  $\Delta P$  which, when superimposed with (a), give the result in (c).

# Modeling of Plastic Zone Under Cyclic Loading

- For proportional plastic flow monotonic eqns used
- Upon load reversal,  $P$  is reduced to  $P - DP$
- Loading parameter replaced by  $DP$
- $s_y$  replaced by  $2s_y$  in monotonic eqns for PZ
- For elastic perfectly plastic solid  $s$  in PZ =  $-s_y$
- For plane stress conditions:  $r_c = (1/\pi)(DK/2s_y)^2$

# Consequences of Reversed Plastic Flow

- Residual plasticity remains even after unloading
- Residual plasticity has implications for VA loading
- Residual stresses self equilibrating
- -ve tip stresses offset by +ve stresses ahead of tip
- Non closing compressive loading induced residual tension
- Cyclic variation induces change in CTOD

# Dugdale Model - Estimates of PZ for Mode I Crack (1960)

- Thin plastic strip of elastic perfectly plastic solid
- Plastic zone loaded by  $s_y$  over  $r_p$
- $r_p = (\pi/8) * (K_I/s_y)^2$
- Similar to Irwin's estimates for plane stress
- CTOD is consequence of necking ahead of crack

## The Dugdale Model

The size of the yield zone ahead of a mode I crack in a thin plate of an elastic-perfectly plastic solid (subject to plane stress deformation) was estimated by Dugdale (1960). If the traction  $\sigma_{yy} = \sigma_x$  were to be applied simultaneously along the length of the strip  $a < |x| < a + r_p$ , it would superimpose a negative stress intensity factor  $K_I''$  on  $K_I'$ , where

$$K_I'' = -\sigma_y \sqrt{\pi(a + r_p)} + 2\sigma_y \sqrt{\frac{(a + r_p)}{\pi}} \sin^{-1}\left(\frac{a}{a + r_p}\right) \quad (1)$$

Since a singular, deviatoric stress state cannot exist at the boundary of the plastic zone,  $K_I' + K_I'' = 0$  (see Hellan, 1984, for further details). Solving for  $r_p$ , one finds that

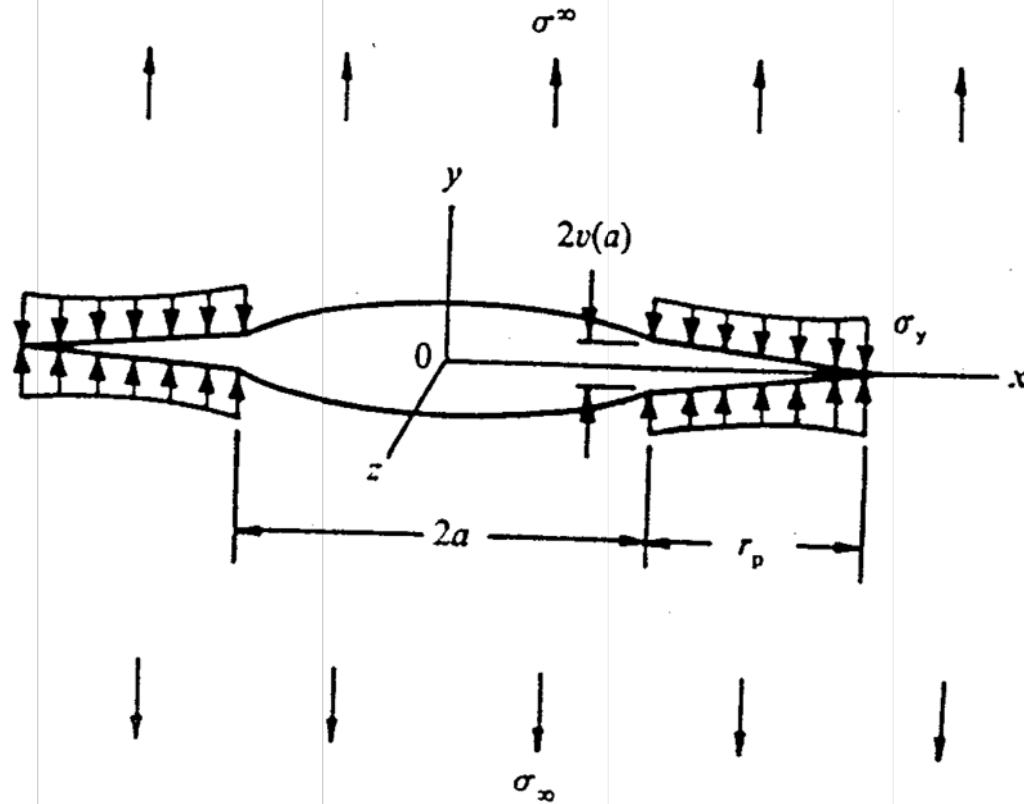
$$\frac{r_p}{a} = \sec\left(\frac{\pi\sigma^\infty}{2\sigma_y}\right) - 1 \quad (2)$$

For  $\sigma^\infty \ll \sigma_y$  and hence for  $r_p \ll a$ , Eq. (1) will asymptotically lead to a plastic zone size

$$r_p = \frac{\pi}{8} \left(\frac{K_I}{\sigma_y}\right)^2$$

This asymptotically exact result due to Dugdale compares well with the Irwin approximation for plane stress.

# A schematic representation of the Dugdale plastic zone model



# Barenblatt Model (1962)

- Analogue to strip yield model for brittle materials
- Consider  $s_{yy}$  = bond rupture strength ( $E/10$ )
- Critical crack size  $f$ (crack-tip cohesive zone)
- Or critical crack size =  $f$ (COD) - (Rice, 1968)

# Elastic-Plastic Fracture Mechanics

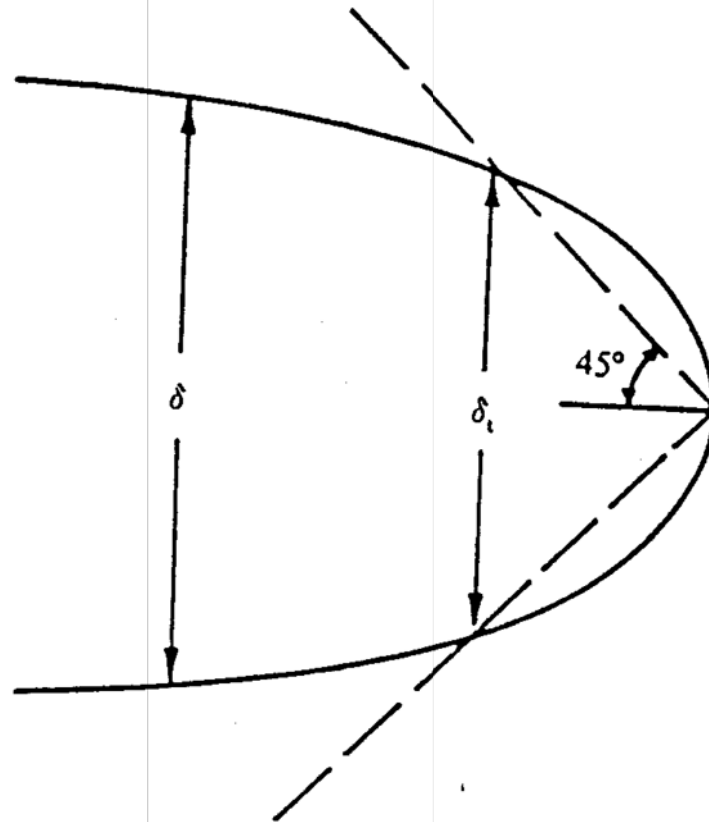
- LEFM valid only for limited plasticity
- EPFM needed for many cases
- CTOD (Wells, 1963)
- J integral (Rice 1968)
- $\Delta J$  of  $\Delta$ CTOD



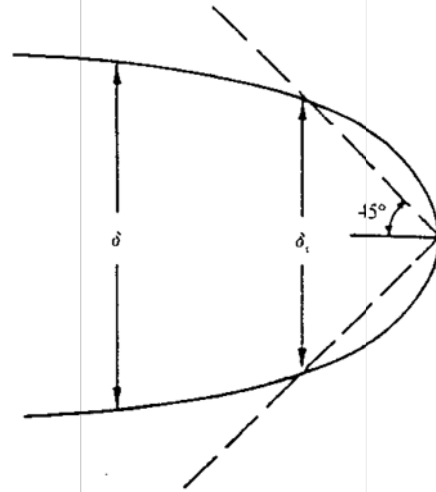
# Crack-Tip Opening Displacement

- CTOD expressions derived from Dugdale model
- More accurate models include effect of hardening
- CTOD definition somewhat arbitrary
- Fracture occurs at critical CTOD
- FCG related to  $\Delta(\text{CTOD})$

## Definition of crack tip opening displacement, $\delta_t$



## Crack Tip Opening Displacement



The definition of  $\delta_t$  is somewhat arbitrary because the distance between the crack faces,  $\delta = \mu_y(x, 0^+) - \mu_y(x, 0^-)$  varies as  $(-x)^{1/(n+1)}$  as the crack tip is approached.

A commonly used operational definition of  $\delta_t$  is based on the distance between two points on the upper and lower crack faces where two  $45^\circ$  lines drawn from the deformed crack tip intercept the crack faces.

$$\delta_t = d_n \frac{J}{\sigma_y}$$

where  $d_n$  is a function of  $\alpha$ ,  $\epsilon_y$  and  $n$ .  $d_n$  ranges in value from about 0.3 to 0.8 as  $n$  is varied from 3 to 13.

## J Integral and Conditions of J-Dominance

- $$J = \int_{\Gamma} \left( \omega dy - T \cdot \frac{\partial u}{\partial x} ds \right)$$

where  $u$  = displacement vector,  $y$  = direction along normal to crack plane,  
 $s$  = arc length,  $T$  = traction vector,  $\omega$  = strain energy density,

$$\sigma_{ij} = \frac{\partial \omega}{\partial \epsilon_{ij}}$$

- For linear elastic and non-linear elastic behavior - J path independent

- Rice (1968) showed that

$$J = G = - \frac{\partial(\text{PE})}{\partial a}$$

- Hutchinson (1983) - J valid when:
  - $J_2$  deformation theory of plasticity gives adequate model of  $\sigma - \epsilon$  behavior
  - Damage and high strain region within HRR field

# Hutchinson-Rice-Rosengreen Singular Fields

- Hutchinson (1968) & Rice and Rosenfield (1968)
- Elastic power law material
- Ramberg-Osgood relationship characterizes matrix behavior
- $J_2$  Deformation Theory
- $J$  is a measure of the intensity of the cracktip fields
- $J = J_c$  when conditions of  $J$  dominance satisfied (ASTM E813)

## Hutchinson-Rice-Rosengreen (HRR) Singular Fields

- Developed for non-linear elastic solids - small strain monotonic deformation
  - Hutchinson (1968)
  - Rice & Rosegreen (1968)
- Near-tip fields given by

$$\sigma_{ij} = \sigma_y \left( \frac{J}{\alpha \sigma_y \varepsilon_y I_n r} \right)^{\frac{1}{n+1}} \tilde{\sigma}_{ij}(\theta, n)$$

$$\varepsilon_{ij} = \alpha \varepsilon_y \left( \frac{J}{\alpha \sigma_y \varepsilon_y I_n r} \right)^{\frac{n}{n+1}} \tilde{\varepsilon}_{ij}(\theta, n)$$

$$u_i = \alpha \varepsilon_y \left( \frac{J}{\alpha \sigma_y \varepsilon_y I_n} \right)^{\frac{n}{n+1}} r^{1/(n+1)} \tilde{u}_i(\theta, n)$$

$\tilde{\sigma}_{ij}(\theta, n)$ ,  $\tilde{\varepsilon}_{ij}(\theta, n)$  and  $\tilde{u}_i(\theta, n)$  are universal functions.

# Conditions for J Dominance

- See review article by Hutchinson (1983)
- Deformation theory of plasticity must be adequate
- This is true for proportional loading (monotonic loading)
- $J_2$  theory not satisfied for elastic-power law plastic solid
- Region of finite strains within process zone
- Finite strain effects significant over  $3 \cdot CTOD$

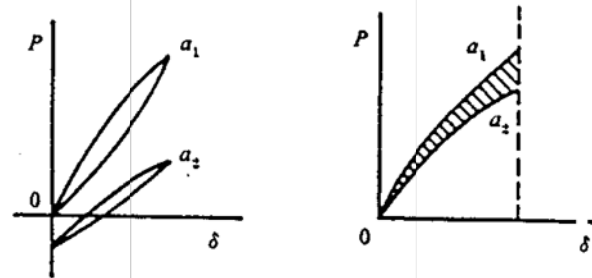
## Condition For J Dominance

**For J-controlled crack growth, Hutchinson & Paris (1979) have suggested that the regime of elastic unloading and nonproportional loading should be confined to well within the zone of J-dominance. In other words,**

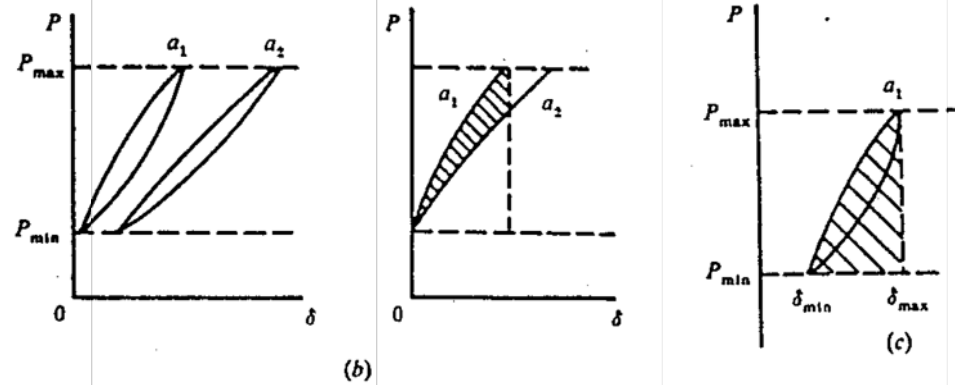
$$\frac{dJ}{da} \gg \frac{J}{R} \text{ and } \Delta a \ll R$$



## Determination of J-integral with stabilized cyclic hysteresis loops.



(a)



(b)

(c)

(a) Hysteresis loops for two different crack lengths in displacement-controlled fatigue and the translation of the rising part of the stabilized hysteresis loop to a common origin. (b) Similar method for load-controlled fatigue with the minimum load being employed as the reference point. (c) Determination of J using a single specimen.

# Fracture Processes and J Dominance

- Region of J dominance must engulf fracture process zone
- HRR solutions hold over 20-25% of PZS in ductile solids
- J dominance specimen dependent for large scale yielding
  - 1% of length of uncracked ligament for CCT panel
  - 7% for deeply notched bend bar (McMeeking & Parks, 1979)
- $R > GS$  for intergranular/transgranular fracture
- $R >$  particle spacing for ductile dimpled fracture